

Theoretical Investigation of Two-Dimensional Superconductivity in Intercalated Graphene Layers

R. A. Jishi^{*} and D. M. Guzman[†]

*Department of Physics and Astronomy,
California State University, Los Angeles, California 90032, USA*

H. M. Alyahyaei[‡]

*Department of Physics and Astronomy,
University of California, Riverside, California 92521, USA*

First-principles calculations of the electronic structure and vibrational modes, in a system of graphene bilayers and trilayers intercalated with alkaline earth atoms, are presented. It is found that, in similarity to the case of superconducting graphite intercalation compounds, the Fermi level is crossed by an s-band derived from the intercalant states, as well as graphitic π -bands. The electron-phonon coupling parameter λ is found to be 0.60 and 0.80, respectively, in calcium intercalated graphene bilayers and trilayers. In superconducting CaC_6 graphite intercalation compound, the calculated value for λ is 0.83. It is concluded that two-dimensional superconductivity is possible in a system of a few graphene layers intercalated with calcium.

PACS numbers: 71.15.Mb, 63.22.Rc, 73.22.Pr, 74.78.-w

Keywords: graphene intercalates, electron-phonon coupling, superconductivity

I. INTRODUCTION

Superconductivity in graphite intercalation compounds (GICs) has been studied extensively for a long period of time [1–10]. A fascinating aspect of these superconducting compounds is that they are usually formed from components neither one of which is a superconductor. In the case of alkali metal GICs the superconducting transition temperature is of the order of 1 K. Within the GIC Brillouin zone (BZ) there exists an alkali-derived spherical band centered at the Γ -point and carbon-derived cylindrical π -bands centered along the HK axis. It was proposed [6, 11, 12] that the presence of both bands is necessary for superconductivity, and that a two-band model is capable of explaining various experimental

^{*}Electronic address: raljish@exchange.calstatela.edu

[†]Electronic address: dguzman4@calstatela.edu

[‡]Electronic address: halya001@ucr.edu

observations relating to the occurrence of superconductivity in alkali-metal GICs and their anisotropic magnetic properties.

It was also found that alkaline earth GICs exhibit superconductivity [7, 8] with a transition temperature of 6.5 K (YbC₆) and 11.5 K (CaC₆). Careful analysis [13–16] of the electron-phonon coupling in CaC₆ reveals that this superconductor is of the conventional BCS type, and that the presence of both Ca-derived and C-derived bands is necessary to explain the relatively high superconducting transition temperature in this compound. It should be pointed out, however, that some controversy still exists regarding the mechanism for superconductivity in CaC₆ GIC, whereas a large Ca isotope effect was measured in CaC₆ GIC [17, 18] indicating a dominant role played by phonons due to Ca atomic vibrations, angle-resolved photoemission spectroscopy (ARPES) measurements on CaC₆ GIC [19], on the other hand, indicate that graphitic high frequency phonon modes are most strongly coupled to electrons.

The recent discovery of stable graphene sheets [20] raises the question of whether two-dimensional superconductivity is possible in intercalated graphene layers. The possibility of intercalating a system consisting of a few graphene layers with intercalants, such as FeCl₃ and Br, has been recently demonstrated [21, 22]. It is clear that a system consisting of many graphene layers, if intercalated, should exhibit properties similar to those of GICs. It is natural to ask whether a system consisting of a few graphene layers, intercalated with alkali, or alkaline earth, metal atoms, will exhibit superconductivity. This problem was discussed recently by Mazin and Balatsky [23], who considered a graphene bilayer intercalated with Ca. They argued, on the basis of similarities with the electronic energy bands of the CaC₆ GIC, that a graphene bilayer intercalated with Ca is probably a superconductor.

In this paper we examine the problem of superconductivity in intercalated graphene layers. In particular, we calculate the energy bands in a graphene bilayer intercalated with the alkaline earth atoms Ca, Sr, and Ba. In all three cases we find that a band, derived from the alkaline earth atoms, crosses the Fermi energy, along with carbon-derived π -bands. We then calculate the electron-phonon coupling strength λ in these systems. The values obtained are reasonably large; in the case of Ca-intercalated bilayer, we find $\lambda = 0.60$, a value to be compared with 0.83 calculated in CaC₆ GIC [24]. Next, the case of a graphene trilayer intercalated with Ca is examined; here we find a value of 0.80 for the electron-phonon coupling strength λ . These calculations indicate that superconductivity is indeed possible in intercalated graphene layers.

In Section II we discuss the computational methods employed in this work. The results of the energy bands, vibrational modes, and electron-phonon coupling calculations are presented and discussed in Section III, and conclusions are given in Section IV.

II. METHODS

The total energy and electronic band structure calculations are carried out using the all-electron, full-potential, linear augmented plane wave (FP-LAPW)

method as implemented in the WIEN2K code [25]. The exchange-correlation potential was calculated using the generalized gradient approximation (GGA) as proposed by Pedrew, Burke, and Ernzerhof (PBE) [26]. The radii of the muffin-tin spheres for the carbon and the alkaline earth atoms were taken as $1.3a_0$ and $2.5a_0$, respectively, where a_0 is the Bohr radius. We set the parameter $R_{MT}K_{max} = 7$, where R_{MT} is the smallest muffin-tin radius and K_{max} is a cutoff wave vector. The valence electrons wave functions inside the muffin-tin spheres are expanded in terms of spherical harmonics up to $l_{max} = 10$, and in terms of plane waves with a wave vector cutoff K_{max} in the interstitial region. The charge density is Fourier expanded up to a maximum wave vector $G_{max} = 13a_0^{-1}$. Convergence of the self-consistent field calculations is attained with a total energy convergence tolerance of 0.01 mRy.

The calculations of the frequencies of the vibrational modes and the electron-phonon coupling parameter are performed using density functional theory (DFT) and ultrasoft pseudopotentials [27, 28]. The electron-phonon coupling strengths are calculated using density functional perturbation theory within the linear response approximation. Ultrasoft pseudopotentials [29] are used for both, the alkaline earth and the carbon atoms. The valence electrons wave function and charge density are expanded in plane waves using 30 Ry and 300 Ry cutoffs, respectively. In the calculation of the vibrational frequencies the electronic integration is carried out using a uniform mesh of $N_k = 8 \times 8$ k-points in the two-dimensional BZ. On the other hand, the electronic density of states, used in computing the electron-phonon coupling parameter λ , is calculated using a finer uniform mesh of 24×24 k-points. The value of λ is obtained by averaging over a uniform mesh of $N_q = 6 \times 6$ phonon momentum q-points.

III. RESULTS AND DISCUSSION

The first issue to decide is this: given a graphene bilayer, will it be energetically favorable for the alkaline earth atoms to be intercalated between the two carbon sheets comprising the bilayer? Or will the energy be lower if the alkaline earth atoms sit on top of the bilayer? If the atoms sit on top of the bilayer, the resulting structure will be denoted by C_6-C_6-A , where A represents an alkaline earth atom. On the other hand, the intercalated structure is denoted by C_6-A-C_6 . In order to decide which structure is more stable, we carried out total energy calculations for both structures. In the case of C_6-C_6-A , the alkaline earth atom is placed over the center of a hexagon in the upper sheet, and its distance to the sheet is varied in order to find the lowest energy. On other hand, in the case of C_6-A-C_6 , the distance between the carbon sheets is fixed at the measured separation in the corresponding graphite intercalation compound [9, 30]. In all the cases considered, corresponding to A=Ca, Sr, and Ba, we found that the energy is lower for the intercalated structure. For the case A=Ca, the intercalated structure C_6-Ca-C_6 is lower in energy than the C_6-C_6-Ca structure by 0.47 eV/Ca atom. In the other two cases, A=Sr and Ba, the energy of the C_6-A-C_6 structure is lower than the C_6-C_6-A structure by 0.43 eV/Sr atom and 0.59 eV/Ba atom. We conclude that it is energetically favorable for the alkaline earth atom to be intercalated in

between the graphene layers.

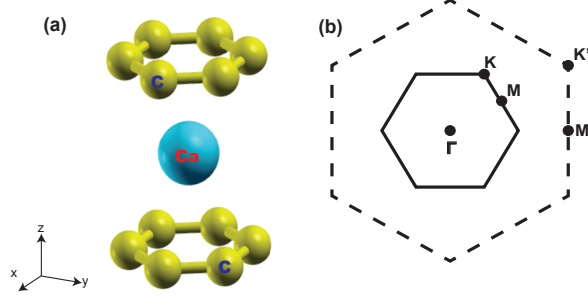


FIG. 1: (Color online) (a) The unit cell of the structure $C_6\text{-A-}C_6$, where A is an alkaline earth atom, is shown. (b) The 2-D Brillouin zone of a graphene layer (or bilayer) is shown as a dashed hexagon, while that for the intercalated graphene bilayer is shown as a solid hexagon. Note that the M and K points in the graphene BZ (here denoted by M' and K') are folded onto points M and Γ , respectively, of the intercalated bilayer BZ.

In the following we focus attention only on the $C_6\text{-A-}C_6$ structure. The separation between the two carbon sheets of the intercalated bilayer is taken as 4.52 Å, 4.94 Å, and 5.25 Å for A=Ca, Sr, and Ba, respectively. The unit cell is a $(\sqrt{3}a \times \sqrt{3}a)$ R 30° supercell of the graphene sheet unit cell; i.e., the in-plane lattice constant of $C_6\text{-A-}C_6$ is $\sqrt{3}$ times the lattice constant in the graphene sheet, and the in-plane lattice vectors of $C_6\text{-A-}C_6$ are rotated by 30° relative to those in the graphene sheet. The intercalated bilayer unit cell, along with the first Brillouin zones of the graphene sheet and the supercell structure are shown in Figure 1. Since the DFT calculations are carried out on a 3-dimensional crystal, the c -lattice constant is taken to be sufficiently large ($c = 14.6$ Å, 15.0 Å, and 15.3 Å for A= Ca, Sr, and Ba, respectively) so that adjacent intercalated bilayers have negligible interactions.

In Figure 2 the calculated energy bands in $C_6\text{-Ca-}C_6$ are plotted in the fatband representation in two different panels, along high symmetry directions in the first BZ. The energy band dispersion curves in both panels are identical, but the width (or the “fatness”) of the band at every k -point in Figure 2a is proportional to the contribution of the Ca atoms to the band eigenstates, while in Figure 2b, the width is proportional to the contribution of the C-derived states. We note that an s-band originating from the Ca states, as well as π -bands originating from C states, cross the Fermi level. This is indicative of the partial charge transfer from the intercalant atoms to the carbon sheets. Indeed, using the Bader’s atoms in molecule (AIM) method[31], we calculated a charge transfer to the carbon sheets of 1.2 electrons/Ca atom, 1.06 electrons/Sr atom, and 0.91 electrons/Ba atom. We also note that in the $(\sqrt{3} \times \sqrt{3})$ R 30° supercell structure, the K-point of the BZ of the graphene sheet is folded onto the Γ point of the BZ of the supercell, as Figure 1 indicates. This explains the observation that the π -bands that are centered at the K-point in graphene are now centered at the Γ -point in the band structure shown in Figure 2. The total and atomic densities of states are plotted

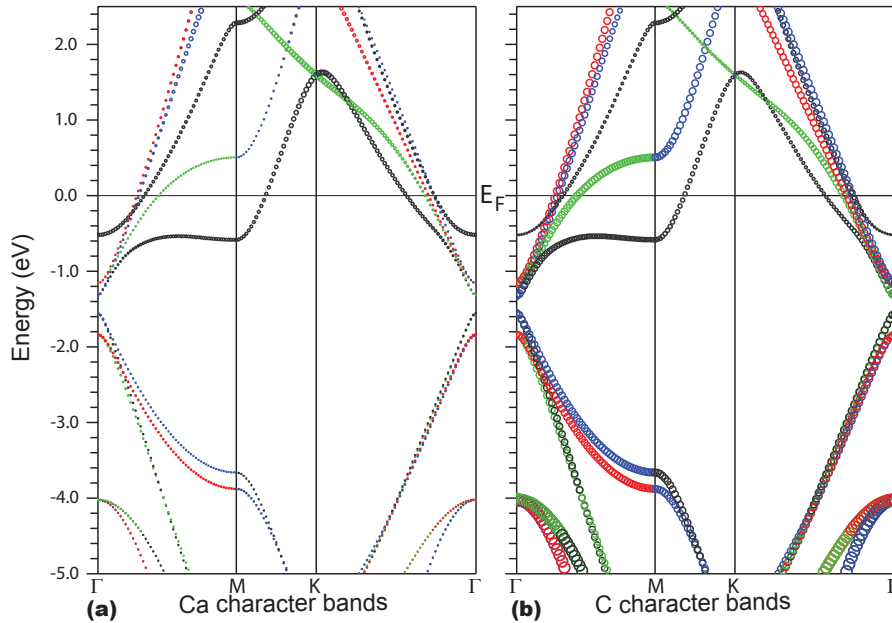


FIG. 2: (Color online) Plot of the energy bands in Ca-intercalated graphene bilayer showing the character of the contributing states. The width of the band at every k-point is proportional to the contribution of (a) Ca-derived states, and (b) C-derived states.

in Figure 3. We note that the total density of states (DOS) consists of the sum of the atomic DOS and the DOS in the interstitial region. The plot in Figure 3 reveals that there is a substantial contribution to the total DOS coming from the interstitial region.

It has been argued previously that for superconductivity in GICs to exist, it is necessary that the charge transfer from the intercalant to carbon sheets be incomplete [11, 12], so that a band derived from the intercalant states must coexist with the π -bands derived from the carbon states. Superconductivity was then explained as resulting from scattering of electrons, through phonon emission or absorption, between the intercalant and the π -bands. More recently, DFT calculations showed that in the CaC_6 GIC, the π -bands alone can not account for the observed superconducting transition temperature, and that the presence of the Ca-derived band is necessary [13–16]. Hence, it is reasonable to expect that superconductivity is possible in the intercalated bilayers under discussion in this work. The superconducting transition temperature will be determined mainly by the strength of the electron-phonon coupling in these systems.

Similar to the case of CaC_6 GICs [24], we calculate the phonon spectrum of the bilayers using DFT as implemented in the PWSCF code [27, 28]. The technical details of the calculation are given in Section II.

For the graphene bilayer intercalated with alkaline earth atoms, the crystal point group is $D_{6h} = D_6 \times i$, where i is the inversion symmetry. The group's character table is given in Table I. The C_6 axis is the z -axis which passes through the alkaline earth atom and is perpendicular to the carbon sheets. The C'_2 axes in the $x - y$ plane are parallel to the lines bisecting the hexagonal sides, and the

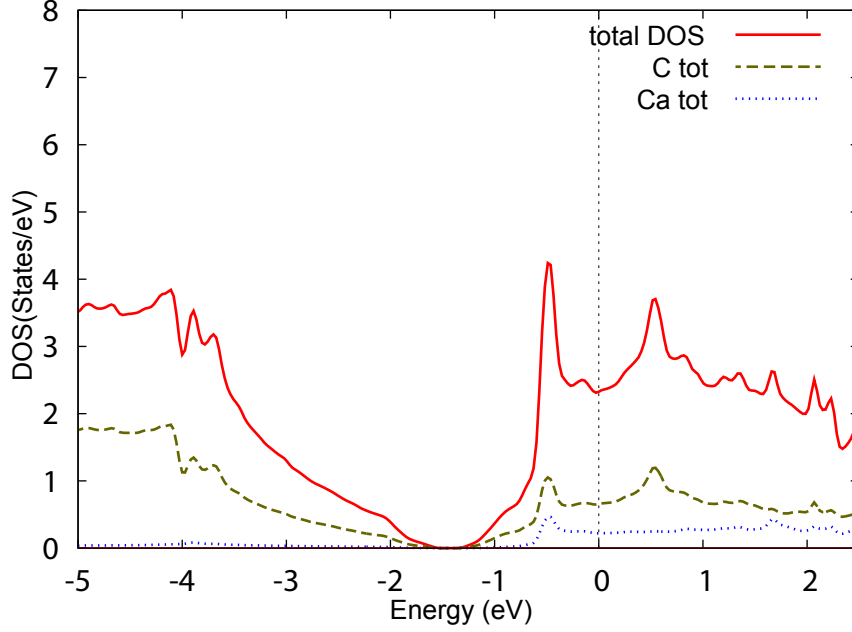


FIG. 3: (Color online) Total and atomic densities of states in Ca-intercalated graphene bilayer. The total density of states is the sum of the atomic densities of states and the density of states in the interstitial region. The Fermi energy is at $E=0$.

C_2'' axes are parallel to the hexagonal diagonals. The $3N = 39$ -dimensional representation, Γ , is given at the bottom of Table I; it is obtained by considering the transformation properties of the 39 cartesian unit vectors placed on the thirteen atoms of the unit cell. From Table I, the 39 phonon modes at the BZ center are decomposed into irreducible representations according to

$$\begin{aligned} \Gamma_{\text{phonon}} = & 2A_{1g} + A_{2g} + 2B_{1g} + B_{2g} + 3E_{1g} + 3E_{2g} \\ & + A_{1u} + B_{1u} + 3A_{2u} + 2B_{2u} + 4E_{1u} + 3E_{2u} \end{aligned}$$

The A_{1g} , E_{1g} , and E_{2g} modes are Raman-active while the A_{2u} and E_{1u} modes are infrared-active. From the representations of x , y , and z , in Table I, the three translational modes are found to transform as A_{2u} and E_{1u} . The Raman and infrared-active modes, with nonzero frequencies, for the three intercalated bilayers considered in this work are given in Table II. The overall trend displayed in Table II is in accordance with what is expected on general grounds. The high frequency modes are lower than the corresponding modes in graphite as a result of the considerable charge transfer which softens the force constants in the carbon sheets. The frequencies of the higher modes increase slightly as the intercalant is changed from Ca to Sr to Ba, a consequence of the concomitant decrease in the amount of charge transfer. The validity of the first-principles calculations of the phonon spectra in the intercalated bilayer may be checked by repeating the same calculation for pristine graphite where the Γ point phonon frequencies are known experimentally. In graphite, the calculated phonon frequencies of the Raman-active modes are $\omega(E_{2g}) = 75 \text{ cm}^{-1}$ and 1557 cm^{-1} , the IR-active modes $\omega(A_{2u}) = 892 \text{ cm}^{-1}$, $\omega(E_{1u}) = 1567 \text{ cm}^{-1}$, and the silent modes $\omega(B_{1g}) = 133$

TABLE I: Character table of $D_{6h} = D_6 \times i$. The character of the 39-dimensional representation, Γ , is given at the bottom.

Axes	Modes	E	C_2	$2C_3$	$2C_6$	$3C'_2$	$3C''_2$	i	iC_2	$2iC_3$	$2iC_6$	$3iC'_2$	$3iC''_2$
z	A_{1g}	1	1	1	1	1	1	1	1	1	1	1	1
	A_{1u}	1	1	1	1	1	1	-1	-1	-1	-1	-1	-1
	A_{2g}	1	1	1	1	-1	-1	1	1	1	1	-1	-1
	A_{2u}	1	1	1	1	-1	-1	-1	-1	-1	-1	1	1
	B_{1g}	1	-1	1	-1	1	-1	1	-1	1	-1	1	-1
	B_{1u}	1	-1	1	-1	1	-1	-1	1	-1	1	-1	1
	B_{2g}	1	-1	1	-1	-1	1	1	-1	1	-1	-1	1
	B_{2u}	1	-1	1	-1	-1	1	-1	1	-1	1	1	-1
(x, y)	E_{1g}	2	-2	-1	1	0	0	2	-2	-1	1	0	0
	E_{1u}	2	-2	-1	1	0	0	-2	2	1	-1	0	0
	E_{2g}	2	2	-1	-1	0	0	2	2	-1	-1	0	0
	E_{2u}	2	2	-1	-1	0	0	-2	-2	1	1	0	0
	Γ	39	-1	0	2	-1	-1	-3	1	0	-2	5	1

cm^{-1} and 889 cm^{-1} . Experimentally, the measured frequencies [32–35] of the Raman-active modes are 42 and 1582 cm^{-1} , the IR-active modes 868 and 1588 cm^{-1} , and the low frequency silent mode 128 cm^{-1} . It is seen that, except for the low frequency mode at 42 cm^{-1} , a good agreement exists between the first-principles calculations and experiment, the error being of the order of a few percent.

TABLE II: Raman- and Infrared-active modes in the intercalated graphene bilayer $\text{C}_6\text{-A-C}_6$ for A=Ca, Sr, and Ba.

		$\omega(\text{irrep})/\text{cm}^{-1}$			
Raman-active modes	Ca	69(E_{1g})	128(A_{1g})	497(E_{1g})	498(E_{2g})
		1132(E_{2g})	1135(E_{1g})	1343(A_{1g})	1410(E_{2g})
	Sr	45(E_{1g})	133(A_{1g})	506(E_{1g})	511(E_{2g})
		1152(E_{2g})	1152(E_{1g})	1355(A_{1g})	1413(E_{2g})
	Ba	38(E_{1g})	115(A_{1g})	506(E_{1g})	509(E_{2g})
		1166(E_{1g})	1167(E_{2g})	1369(A_{1g})	1454(E_{2g})
Infrared-active modes	Ca	115(E_{1u})	293(A_{2u})	481(E_{1u})	1130(E_{1u})
		1349(A_{2u})			
	Sr	72(E_{1u})	207(A_{2u})	498(E_{1u})	1151(E_{1u})
		1358(A_{2u})			
	Ba	67(E_{1u})	153(A_{2u})	503(E_{1u})	1165(E_{1u})
		1370(A_{2u})			

We calculated the electron-phonon coupling parameter λ in $\text{C}_6\text{-A-C}_6$ and found it to be 0.60, 0.42, and 0.35, for A=Ca, Sr, and Ba respectively. The transition

superconducting temperature T_c of the the studied compounds was estimated using the Allen-Dynes [36] formula with values for the phonon logarithmic average frequencies $\omega_{log} = 29.4$ meV, 30.1 meV and 28.4 meV for $C_6\text{-Ca-C}_6$, $C_6\text{-Sr-C}_6$, and $C_6\text{-Ba-C}_6$ respectively. Using a screened Coulomb pseudopotential $\mu^* = 0.14$ we obtained $T_c = 4.7$ K for Ca-intercalated graphene bilayer, 0.95 K for Sr-intercalated bilayer, and 0.1 K for Ba-intercalated bilayer. This is to be compared with the value $\lambda = 0.83$ calculated previously in CaC_6 GIC [24] for which T_c is 11.5 K.

In order to study the effect of increasing number of graphene layers on the superconducting properties, we consider a trilayer of graphene intercalated with calcium: Ca_2C_{18} . The unit cell is again given by $(\sqrt{3}a \times \sqrt{3}a) R 30^\circ$ supercell of the graphene sheet unit cell. In conformity with the corresponding situation in stage-1 GICs, the carbon layer stacking is taken to be AAA. Between the first and second carbon layers, Ca atoms sit above hexagon centers, occupying positions labeled α . Between the second and third carbon layers, Ca atoms sit above hexagon centers, occupying positions labeled β . The two intercalant layers are not directly above each other; the second intercalant layer is shifted, relative to the first intercalant layer, by a lattice vector of the graphene sheet, as illustrated in Figure 4. Thus, the stacking in the intercalated graphene trilayer is $A\alpha A\beta A$. The crystal point group is reduced, from D_{6h} in the case of intercalated bilayer, to D_{3d} in the case of the intercalated trilayer. The phonon modes, at the BZ center, are decomposed into the following irreducible representations

$$\Gamma_{phonon} = 5A_{1g} + 5A_{2g} + 4A_{1u} + 6A_{2u} + 10E_g + 10E_u.$$

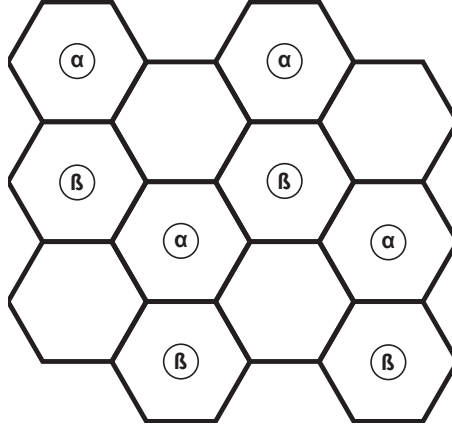


FIG. 4: (Color online) Projection, on the plane of the first carbon layer, of the intercalated graphene trilayer. In the first intercalant layer, Ca atoms occupy positions denoted by α , whereas Ca atoms occupy positions β in the second intercalant layer.

The A_{1g} and E_g modes are Raman-active, whereas A_{2u} and E_u modes are IR-active. The three zero-frequency translational modes transform according to A_{2u} and E_u irreducible representations. In Table III we report the Raman- and IR-active modes with non-vanishing frequencies. The electron-phonon coupling parameter in calcium intercalated graphene trilayer is calculated to be $\lambda = 0.80$.

TABLE III: Raman- and Infrared-active modes, with nonzero frequencies, in Ca-intercalated graphene trilayer.

	$\omega(\text{irrep})/\text{cm}^{-1}$				
Raman-active modes	40(E_g)	83(A_{1g})	102(E_g)	249(A_{1g})	411(E_g)
	489(E_g)	492(E_g)	1112(E_g)	1131(E_g)	1133(E_g)
	1326(A_{1g})	1344(A_{1g})	1362(E_g)	1382(A_{1g})	1440(E_g)
Infrared-active modes	65(E_u)	106(E_u)	159(A_{2u})	285(A_{2u})	425(E_u)
	489(E_u)	492(E_u)	1111(E_u)	1132(E_u)	1133(E_u)
	1326(A_{2u})	1344(A_{2u})	1362(E_u)	1401(A_{2u})	

The calculated phonon logarithmic average frequency is 25.2 meV, and the calculated superconducting transition temperature is 10.1 K.

IV. CONCLUSIONS

We have carried out energy band calculations, within density functional theory, of a graphene bilayer intercalated with the alkaline earth atoms A=Ca, Sr, or Ba, and a graphene trilayer intercalated with Ca. In all cases, we found that the Fermi level is crossed by an s-band derived from the intercalant states as well as graphitic π -bands, in similarity to the situation in superconducting graphite intercalation compounds. The electron-phonon coupling parameter λ is largest for the case A=Ca, where $\lambda = 0.60$ for the intercalated bilayer, and $\lambda = 0.80$ for the intercalated trilayer; this is to be compared with $\lambda = 0.83$ for the case of CaC₆ GIC. On the basis of these calculations, we expect that graphene bilayers and trilayers, intercalated with calcium, will exhibit two-dimensional superconductivity. Even though there are no experimental reports yet of the successful intercalation of calcium into a system of a few graphene layers, it is hoped that this work will stimulate further experimental studies in that direction.

Acknowledgments

One of the authors (R. A. J) gratefully acknowledges support by NSF under grant No. HRD-0932421. D. M. G. would also like to thank the CSULA-MORE programs and CSU-LSAMP fellowship for their support under grant No. HRD-1026102-518741.

-
- [1] N. B. Hannay, T. H. Geballe, B. T. Matthias, K. Andres, P. Schmidt and D. MacNair, Superconductivity in Graphitic Compounds, Phys. Rev. Lett. **14**, 255 (1965).
 - [2] Y. Koike and S. Tanuma, Superconductivity in the graphite-potassium intercalation compound C₈K, J. Phys. Chem. Solids **41**, 1111 (1980).

- [3] M. Kobayashi and I. Tsujikawa, Susceptibility Fluctuation in Superconducting First Stage Potassium Graphite Intercalation Compound near Transition Temperature, *J. Phys. Soc. Jpn.* **46**, 1945 (1979).
- [4] M. G. Alexander, D. P. Goshorn, D. Gurard, P. Lagrange, M. El Makrini and D. G. Onn, Superconductivity of the graphite intercalation compounds KHgC8 and RbHgC8: Evidence from specific heat, *Solid State Commun.* **38**, 103 (1981).
- [5] Y. Iye and S. Tanuma, Superconductivity of graphite intercalation compounds with alkali-metal amalgams, *Phys. Rev. B* **25**, 4583 (1982).
- [6] R. Al Jishi, Model for superconductivity in graphite intercalation compounds, *Phys. Rev. B* **28**, 112 (1983).
- [7] N. Emery, C. Herold, J.-F. Mareche, and P. Lagrange, Synthesis and superconducting properties of CaC_6 , *Sci. Technol. Adv. Mater.* **9**, 044102 (2008).
- [8] T. E. Weller, M. Ellerby, S. S. Saxena, R. P. Smith and N. T. Skipper, Superconductivity in the Ytterbium Intercalated Graphite Compounds C_6Yb and C_6Ca , *Nature Physics* **1**, 39 (2005).
- [9] N. Emery, C. Herold, M. dAstuto, V. Garcia, Ch. Bellin, J. F. Mareche, P. Lagrange, and G. Louprias, Superconductivity of Bulk CaC_6 , *Physical Review Letters* **95**, 087003 (2005).
- [10] A. C. Walters, C. A. Howard, M. H. Upton, M. P. M. Dean, A. Alatas, B. M. Leu, M. Ellerby, D. F. McMorrow, J. P. Hill, M. Calandra, F. Mauri, Comparative study of the phonons in non-superconducting BaC_6 and superconducting CaC_6 using inelastic x-ray scattering, *arXiv:1102.4767v1 [cond-mat.supr-con]*, (2011)
- [11] R. A. Jishi, M. S. Dresselhaus, and A. Chaiken, Theory of the upper critical field in graphite intercalation compounds, *Phys. Rev. B* **44**, 10 248 (1991)
- [12] R. A. Jishi and M. S. Dresselhaus, Superconductivity in graphite intercalation compounds, *Phys. Rev. B* **45**, 12465 (1992).
- [13] M. Calandra and F. Mauri, Theoretical Explanation of Superconductivity in C_6Ca , *Physical Review Letters* **95**, 237002 (2005).
- [14] M. Calandra and F. Mauri, Possibility of superconductivity in graphite intercalated with alkaline earths investigated with density functional theory, *Physical Review B* **74**, 094507 (2006).
- [15] A. Sanna, G. Profeta, A. Floris, A. Marini, E. K. U. Gross and S. Massidda, Anisotropic gap of superconducting CaC_6 : A first-principles density functional calculation, *Physical Review B* **75**, 020511 (2007).
- [16] M. Calandra, G. Profeta, F. Mauri, Adiabatic and nonadiabatic phonon dispersion in a Wannier function approach, *Physical Review B* **82**, 165111 (2010).
- [17] D. G. Hinks, D. Rosenmann, H. Claus, M. S. Bailey, and J. D. Jorgensen, Large Ca isotope effect in the CaC_6 superconductor, *Phys. Rev. B* **75**, 014509 (2007).
- [18] R. S. Gonnelli, D. Daghero, D. Delaude, M. Tortello, G. A. Ummarino, V. A. Stepanov, J. S. Kim, R. K. Kremer, A. Sanna, G. Profeta, and S. Massidda, Evidence for Gap Anisotropy in CaC_6 from Directional Point-Contact Spectroscopy, *Phys. Rev. Lett.* **100**, 207004 (2008).
- [19] T. Valla, J. Camacho, Z. H. Pan, A. V. Fedorov, A. C. Walters, C. A. Howard, and M. Ellerby, Anisotropic Electron-Phonon Coupling and Dynamical Nesting on the Graphene Sheets in Superconducting CaC_6 using Angle-Resolved Photoemission

- Spectroscopy, Phys. Rev. Lett. **102**, 107007 (2009).
- [20] K. S. Novoselov, A. K. Geim, S. V. Morozov, D. Jiang, Y. Zhang, S. V. Dubonos, I. V. Grigorieva, and A. A. Firsov, Electric Field Effect in Atomically Thin Carbon Films, Science **306**, 5696 (2004).
 - [21] N. Jung, N. Kim, S. Jockusch, N. J. Turro, P. Kim and L. Brus, Charge Transfer Chemical Doping of Few Layer Graphenes: Charge Distribution and Band Gap Formation, Nano Letters **9**, 4133 (2009)
 - [22] N. Jung, A. C. Crowther, N. Kim, P. Kim and L. Brus, Raman Enhancement on Graphene: Adsorbed and Intercalated Molecular Species, ACS Nano **4**, 7005 (2010).
 - [23] I.I. Mazin, A.V. Balatsky, Superconductivity in Ca-intercalated bilayer graphene, arXiv:0803.3765v1 [cond-mat.supr-con], (2008).
 - [24] M. Calandra and F. Mauri, Origin of superconductivity of CaC_6 and of other intercalated graphites, Physica Status Solidi(b) **243**, 3458 (2006).
 - [25] P. Blaha, K. Schwarz, G. K. H. Madsen, D. Kvasnicka, and J. Luitz, WIEN2K: *An Augmented Plane Wave+Local Orbitals Program for Calculating Crystal Properties* (Technische Universität, Wien, Austria, 2001).
<http://www.wien2k.at>.
 - [26] J. P. Perdew, K. Burke, and M. Ernzerhof, Generalized Gradient Approximation Made Simple, Phys. Rev. Lett. **77**, 2865 (1996).
 - [27] S. Scandolo et al., First-principles codes for computational crystallography in the Quantum-ESPRESSO package, Zeit. Krist. **220**, 574 (2005).
 - [28] P. Giannozzi et al., QUANTUM ESPRESSO: a modular and open-source software project for quantum simulations of materials, J. Phys. Condens. Matter **21**, 395502 (2009).
<http://www.quantum-espresso.org/>
 - [29] D. Vanderbilt, Soft self-consistent pseudopotentials in a generalized eigenvalue formalism, Phys. Rev. B **41**, 7892 (1990).
 - [30] D. Guerard, M. Chaabouni, P. Lagrange, M. El Makrini, and A. Herold, Insertion de metaux alcalino-terreux dans le graphite, Carbon **18**, 257 (1980).
 - [31] R. Bader, *Atoms in Molecules: A Quantum Theory* (Clarendon Press, Oxford, UK, 1994).
 - [32] F. Tunista and J. L. Koenig, Raman Spectrum of Graphite, J. Chem. Phys. **53**, 1126 (1970).
 - [33] L. J. Brillson, E. Burstein, A. A. Maradudin, and T. Stark, in *Proceedings of the International Conference on the Physics of Semimetals and Narrow Gap Semiconductors*, Dallas, Texas, 1970, edited by E. L. Carter and R. T. Bate (Pergamon, New York, 1971), p. 187.
 - [34] R. J. Nemanich, G. Lucovsky, and S. A. Solin, Infrared active optical vibrations of graphite, Solid State Commun. **23**, 117 (1977).
 - [35] R. Nicklow, N. Wakabayashi, and H. G. Smith, Lattice Dynamics of Pyrolytic Graphite, Phys. Rev. B **5**, 4951 (1972).
 - [36] P. B. Allen and R. C. Dynes, Transition temperature of strong-coupled superconductors reanalyzed, Phys. Rev. B **12**, 905 (1975).